

Last time: we can construct ^{(n+1)-node} Gaussian quadratures if can find $n+1$ degree poly q

- i) orthog. to all P_n , in our inner prod (\cdot, \cdot) . $q_{n+1}(x) = \prod_{j=0}^n (x-x_j)$
- ii) with all roots $x_j \in [a,b]$, which give the nodes

Lemma 9.15 \exists unique seq. $(q_n)_{n=0}^\infty$ w/ $q_0=1$ & $q_n(x) = x^n + p(x)$, $p \in P_n$ which are mutually orthog. $q_n \perp q_m$, $n \neq m$, and $\text{Span}\{q_0, \dots, q_n\} = P_n$

Pf: $1, x, x^2, \dots$ are Lin. Indep., so Gram-Schmidt unique:

$$q_0 = 1$$

$$q_1 = x - \frac{(x, q_0)}{(q_0, q_0)}$$

$$q_2 = x^2 - \frac{(x^2, q_0)}{(q_0, q_0)} - \frac{(x^2, q_1)}{(q_1, q_1)}$$

$$\vdots$$

$$q_n = x^n - \sum_{j=0}^{n-1} \frac{(x^n, q_j)}{(q_j, q_j)} q_j$$

n L.I. vecs. in P_n ^{must} span it.

'Legendre' poly's (but: not ^{standard} normalization): unique ^(unweighted) orthog. poly's on $[-1, 1]$

Lemma 9.16 q_n has n simple zeros in (a, b) .

Pf: $\forall n \geq 1, q_n \perp q_0$ i.e. $\int q_n = 0$ so q_n has ≥ 1 zeros x_1, \dots, x_m in (a, b) .

Suppose $m < n$, then $r_m = \prod_{j=1}^m (x-x_j) \in P_{n-1}$ so is $\perp q_n$
 But $\int r_m q_n \neq 0$ since $r_m q_n$ has fixed sign, not $\equiv 0$. } contradicts. $\Rightarrow m=n$

In practice, how do we compute nodes $\{x_j\}_{j=0}^n$? They are sigls of $\begin{bmatrix} 0 & \beta_1 & & & \\ \beta_1 & 0 & \beta_2 & & \\ & \beta_2 & 0 & \beta_3 & \\ & & \beta_3 & 0 & \ddots \\ & & & & \ddots & \beta_{n-1} \end{bmatrix}$ ^(tridiagonal) ^{since β} ^{3-term} ^{recurrence}

$$\beta_n := \frac{1}{2}(1 - (2n)^{-2})^{-1/2}$$

& weights w_j come from l^2 component of eigenvectors. [Wilf]
 See gauss.m code, Golub-Welsch scheme.
 We could prove but won't. (beautiful)

Claim: $2n+1$ is highest possible degree of $(n+1)$ -node quadrature.

pf: $p = \prod_{j=0}^n (x-x_j)^2 \in P_{2n+2}$ has $Q_n(p) = 0$ but $Q(p) > 0$.

Thm: Gaussian weights non-neg: pf $l_k(x_j) = \delta_{jk}$ so $l_k^2(x_j) = \delta_{jk}$ also.

so $0 < \int l_k^2(x) dx = \sum_{j=0}^n w_j l_k^2(x_j) = w_k$

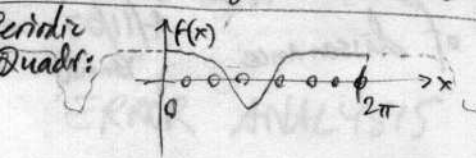
Cor: Gauss. quad ^{nonneg} convergent.

^{exact since} $l_k^2 \in P_{2n}$

There are error bounds for Gauss. quadr., won't do].

This all generalizes to weighted quadrature $Q_n(f) := \int_{-1}^1 f(x) u(x) dx$, in which case

② 10/16/0



weight (we did $u \equiv 1$)
inner prod $(f, g) = \int_{-1}^1 f(x)g(x)u(x)dx$.

Useful for f with singularities
eg. $f(x) = |g(x)| \ln|x|$
smooth \uparrow $u(x)$

Periodic trapezoid rule: $Q_n(f) = \frac{2\pi}{n} \sum_{j=1}^n f\left(\frac{2\pi j}{n}\right)$

w_j all equal $\leftarrow x_i$ equally spaced.

could derive via interpolation of f by trigonometric polys, ie Fourier series truncated at term $\frac{N}{2} \dots$ later.

For now: some error bounds, error $E_n(f) := Q_n(f) - Q(f)$, a number.

Thm (9.27)

Let $f: \mathbb{R} \rightarrow \mathbb{R}$ be 2π -periodic & C^{2m+1} , $m \in \mathbb{N}$.

then $|E_n(f)| \leq C_m \int_a^{2\pi} |f^{(2m+1)}(x)| dx \cdot \frac{1}{n^{2m+1}}$

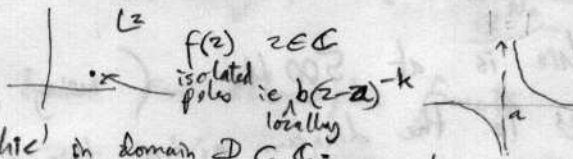
convergence

- means $g \in C^5$, ie $f^{(5)}$ cont. but $f^{(6)}$ discont., quadr. error is $O(n^{-5})$
- So, smoother f gives higher-order algebraic convergence.
- Euler-Maclaurin expansion.

5th order worth effort
($n = \#$ func eval)
 $n \rightarrow 10n$ gets you 10^5 , ie 5 extra digits need $10^3 n$

What is the smoothest type of func? analytic.

Review complex analysis:

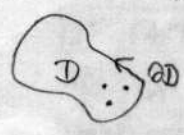


residue.
 $k=1$: simple pole = $\frac{b}{z-a}$

$f(z)$ 'holomorphic' in domain $D \subset \mathbb{C}$; no poles inside D , Taylor series converges in some disc. eg. $\frac{1}{1+25x^2}$ has simple poles at $\pm i/5$.

Residue thm: if f holomorphic in D apart from finite # poles,

then $\int_{\partial D} f(z) dz = 2\pi i \sum_{\text{simple poles}} (\text{residue of each simple pole})$.



- Why? i) Cauchy f holom $\Rightarrow \int_{\partial D} f(z) dz = 0$
ii) small loops around pole cancel unless $f(z)$ goes CW once (since dz goes CW once)

Skip to strip theorem & proof.

Note $\cot \frac{\eta z}{2} \rightarrow \begin{cases} -1 & \text{for } \text{Im} z > 0 \\ +1 & < 0 \end{cases}$ as $n \rightarrow \infty$, exponentially fast. show pic.

Note: Gaussian & Newton-Cotes (worse) & periodic trap. example of 'spectral methods', ie exp. conv. error $O(k^{-n})$

You may also take derivatives of (trig) polys to get formulae for derivatives f', f'' etc, then solve ODEs (or PDEs) with spectral accuracy. Tie e^{-n}

Math 116. [Lecture 6 from 2006]

① 1/24/06
Barnett

ERROR ANALYSIS OF INTEGRATION OF PERIODIC FUNCS.

Why is crude equal-weight equally-spaced quadrature $\int_0^{2\pi} g(x) dx \approx \frac{2\pi}{N} \sum_{j=1}^N g(\frac{2\pi j}{N})$ so good?
 $\sum w_j$ all equal.

ANALYTIC CASE (§9.4, Kress, "Numerical Analysis").

Thm. Let $g: \mathbb{R} \rightarrow \mathbb{R}$ be analytic & 2π -periodic. Then there exists a strip

$D = \mathbb{R} \times (-a, a) \subset \mathbb{C}$ with $a > 0$ s.t. g can be extended to a holomorphic and 2π -periodic bounded function $g: D \rightarrow \mathbb{C}$.

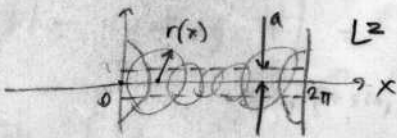
The error for above quadrature rule is bounded by

$$|R_N[g]| \leq \frac{4\pi M}{e^{Na} - 1}$$

where M is a bound for holomorphic function g on D .

Remark: • this proves exponential convergence of errors $O(e^{-aN})$, a vertical disk to nearest pole M f.
 • you'll test in HW.

Proof: **1st PART**



Analytic \Rightarrow at each $x \in \mathbb{R}$, Taylor expansion converges in some open disk radius $r(x) > 0$.

This provides a 2π -periodic holomorphic extension of g .
 \hookrightarrow since x & $x+2\pi$ have same Taylor expansion.

Can cover $[0, 2\pi]$ with finite # of such disks.

a can be chosen to be any width $<$ minimum $r(x)$.

g is then bounded on the strip D .

2nd PART

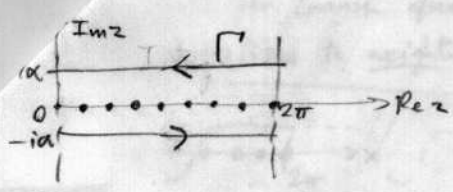
Consider $\frac{1}{z} \cot(z)$, which has residuals (pole strengths)

of 1 at $z_j = \pi j$, $j \in \mathbb{Z}$

(since $\frac{d}{dz} \tan z = \frac{1}{\cos^2 z}$)

Then $g(z) \cot(\frac{N}{2}z)$ has residuals $\frac{2}{N} g(\frac{2\pi j}{N})$

at points $z_j = \frac{2\pi j}{N}$



Residue thm gives, for $\alpha < a$,

$$\int_{\Gamma} g(z) \cot\left(\frac{Nz}{2}\right) dz = 2\pi i \sum \text{residuals}$$

$$= \frac{4\pi i}{N} \sum_{j=1}^N g\left(\frac{2\pi j}{N}\right) \quad (*)$$

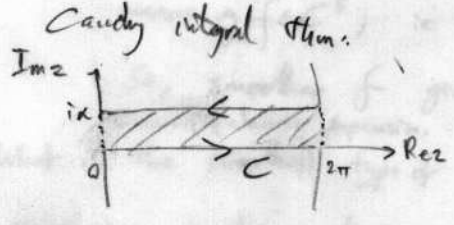
Note periodic \rightarrow doesn't need to close.

Schwarz reflection principle: $\begin{cases} g \text{ real on } \mathbb{R} \text{ so } g(\bar{z}) = \overline{g(z)} \\ \text{ie imaginary part is antisymmetric in } \text{Im } z. \end{cases}$

\Rightarrow LHS integral becomes $-i \int_{ia}^{ia+2\pi} 2 \text{Im } g(z) \cot\left(\frac{Nz}{2}\right) dz = 2i \text{Re} \int_{ia}^{ia+2\pi} i g(z) \cot\left(\frac{Nz}{2}\right) dz$

Using (*), $\text{Re} \int_{ia}^{ia+2\pi} i \cot\left(\frac{Nz}{2}\right) g(z) dz = \frac{2\pi}{N} \sum_{j=1}^N g\left(\frac{2\pi j}{N}\right)$

our quadrature rule!



$\oint_C g(z) dz = 0$ since analytic in D true integral

so $\text{Re} \int_{ia}^{ia+2\pi} g(z) dz = \int_0^{2\pi} g(x) dx$

\Rightarrow error $R_N[g] = \text{Re} \int_{ia}^{ia+2\pi} \left[1 - i \cot\left(\frac{Nz}{2}\right)\right] g(z) dz$

bounded? $y = \frac{ny}{2}$ $|g(z)|$ on $x+iy$ is bounded by M .

$\left|1 - i \cot y \frac{e^{2iy} + 1}{2}\right| = \left|1 + \frac{1 + e^{-2iy}}{1 - e^{-2iy}}\right| = \left|\frac{2}{e^{-2iy} - 1}\right| \leq \frac{2}{e^{N\alpha} - 1}$

Euler for sin, cos. use $|e^{-2iy}| = e^{-2\text{Im } y} = e^{-N\alpha}$

Take limit $\alpha \rightarrow a$. QED.

Remark: $\frac{1}{\pi N} \text{Im} \cot \frac{Nz}{2}$ is just an approximation to double layer potential placed along the Re axis



There also exist Euler-Maclaurin theorems for C^{2m+1} FUNCTIONS:

Thm: Let $g \in C^{2m+1}$ be 2π -periodic, for some $m \geq 1$.

Then $|R_N[g]| \leq \frac{C}{N^{2m+1}} \int_0^{2\pi} |g^{(2m+1)}(x)| dx$ where $C = 2 \sum_{k=1}^{\infty} \frac{1}{k^{2m+1}}$

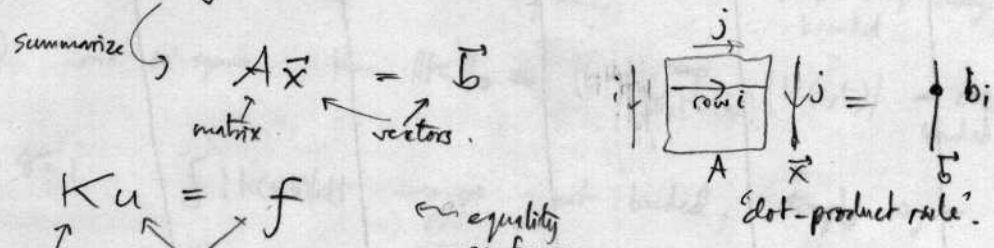
Proof requires Bernoulli poly's (see Kress §9.4).

Smoother $g \Rightarrow$ higher-order convergence

Integral Eqns.

Given Interval (a,b) } Seek function $u(t)$ solving $\int_a^b k(s,t) u(t) dt = f(s)$ for $s \in (a,b)$
 func f
 func of 2 var k
 'Fredholm eqn of 1st kind'.
 kernel func on $[a,b]^2$ (the square)
 right-hand side.

Compare linear sys. of eqns. $\sum_j a_{ij} x_j = b_i$ $i=1 \dots n$



Summarize IE by $Ku = f$
 K lin. operator. u func. f equality as func.
 defined by $(Ku)(s) = \int_a^b k(s,t) u(t) dt$

eg. Fourier trans $k(s,t) = e^{ist}$ on \mathbb{R} .
 convolution: blurring of image $k(s,t) = k(s-t)$ on \mathbb{R} .
 Lin ops like square matrices: i) they have eigenvalues & eigfunc, even SVD.
 ii) you can apply them repeatedly, eg $(K^2 u)(s) = \int \int k(s,t) k(t,r) u(r) dr dt = \int k(s,r) u(r) dr$ (like convolution: $\int k(s,t) k(t,r) = \int k(s,r)$)
 Fredholm eqn 2nd kind: $u - Ku = f$ is $(I - K)u = f$
 sometimes nonzero const. here.
 We will use these to solve PDEs. First we solve IEs numerically. This will involve quadrature, linear systems, func. analysis, matrices (finite-dim K brings many twists).

eg. $k(s,t) = \begin{cases} 0 & s \leq t \\ 1 & s > t \end{cases}$



'lower-triangular', or 'causal' op.

$(Ku)(s) = \int_a^s u(t) dt = f(s)$ unique solution is $u = f'$
 indep. integral exists iff $f(0) = 0$ & f diff'ble

This example of Volterra IE: $k=0$ for $s < t$ (above diag). Easy to solve (equiv. to ODE) won't concern us.
 Fredholm have stuff on both sides of diag

eg. $\int_0^1 s^2 t u(t) dt = \frac{1}{3} s^2$ easily solve, bring out s^2 : $\int_0^1 t u(t) dt = \frac{1}{3}$
 $k(s,t)$
 K is rank-1 since $(Ku)(s) = c s^2$ highly nonunique, typical of 1st kind.
 $u(t) = t + (\text{any func } \perp t)$

Bounded operators:

$\|K\| = \sup_{\|u\|=1} \|Ku\|$ for norms of your choice on functions, eg. $L^2(a,b)$ norm (sup norm) $\|u\|_\infty := \max_{s \in [a,b]} |u(s)|$

What is op. norm in terms of kernel?

$$|(Ku)(s)| = \left| \int_a^b k(s,t) u(t) dt \right| \leq \int_a^b |k(s,t)| dt \quad \text{if } \|u\|_\infty = 1.$$

with equality as $u(t) \rightarrow \text{sgn}(k(s,t))$
(see Thm 12.5 proof)

$$\|K\|_\infty = \sup_{s \in [a,b]} |(Ku)(s)| = \sup_{s \in [a,b]} \int_a^b |k(s,t)| dt$$

bigger row-integral of abs val of kernel

Say $b-a < \infty$ and $k \in C([a,b]^2)$ cont. on square, then $\|K\|_\infty \leq (b-a) \sup_{s,t \in [a,b]} |k(s,t)| < \infty$ bounded

But $k(s,t) = \frac{1}{|s-t|^\alpha}$, $\alpha=1$: $\int |k(s,t)| dt \rightarrow \infty$ not bounded, strongly singular.

$0 < \alpha < 1$ integrable \rightarrow bounded, 'weakly-singular' kernel. (blows up, discontinuous at $s=t$)

May instead use L^2 norm: $\|u\|_2 := \sqrt{\int_a^b |u(s)|^2 ds}$, Eg. $|(Ku)(s)| \leq \int_a^b |k(s,t)u(t)| dt \leq \sqrt{\int_a^b |k(s,t)|^2 dt} \sqrt{\int_a^b |u(t)|^2 dt}$

Numerical solution method: Nyström (1930) $(Ku)(s) = \int_a^b k(s,t) u(t) dt = f(s) \quad s \in [a,b]$

is upper bound for $\|K\|_2$.
if $< \infty$, called Hilbert-Schmidt kernel

2nd kind

$$u(s) - \int_a^b k(s,t) u(t) dt = f(s) \quad s \in [a,b]$$

quadr. Q_n

Approx u by u_n which solves:

$$u_n(s) - \sum_{j=0}^n w_j k(s, t_j) u_n(t_j) = f(s) \quad (*)$$

$(K_n u_n)(s)$

Then values at nodes $u_n(s_i) = u_n(s_i)$ sat. the lin. sys: where K_n is rank- n approx to op. K .
its range is $\text{span}\{k(\cdot, t_j)\}$

$$u_i^{(n)} - \sum_{j=0}^n w_j k(s_i, s_j) u_j^{(n)} = f(s_i) \quad (LS)$$

ie $(I - A) \vec{u}^{(n)} = \vec{f}$

\vec{f} vector of RHS at nodes.

So you've solved for u at nodes - how get full func $u_n(s)$?

Thm (12.11) If $\{u_i^{(n)}\}_{i=0}^n$ is soln. to (LS) then $u_n(s) = f(s) + \sum_{j=0}^n w_j k(s, t_j) u_j^{(n)}$ solves $(*)$

pf: $u_n(s_j) = u_j^{(n)} \quad \forall j$ by construction of (LS) soln. (N)

use to sub. for $u_j^{(n)}$ in (N) gives $(*)$ QED. Subtle!

$(*)$ expresses u_n as $f + \text{span}\{\text{column slices of kernel at node } k(\cdot, t_j)\}$, ie. interpolation basis.

(LS) is equiv. of Vandermonde sys to require interpolant agrees at nodes, (N) reconstructs interpolant from node values.